

# Gap probabilities for the Bures–Hall ensemble and the Cauchy–Laguerre two-matrix model

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Baylor Analysis Fest

With L. Wei, TTU

[Hall 1998]

*What statistical ensemble corresponds to minimal prior knowledge about a quantum state?*

- 1 What is a typical quantum state?
  - The setting for quantum mechanics
  - The Fubini-Study metric for projective Hilbert spaces
  - Density matrices
  - The Bures-Hall measure
  - The joint probability density function for eigenvalues of the density matrix
- 2 A statistic for the Bures-Hall ensemble: the gap probability
  - Pfaffian point processes
  - Cauchy-Laguerre Matrix model
  - Bi-orthogonal system of polynomials
  - Integrability and the dynamical system

- Complex Hermitian  $n \times n$  matrix  $\rho = \rho^\dagger$
- Unit trace  $\text{Tr}\rho = 1$
- Positive definite  $\rho > 0$ , i.e.  $\forall |\psi\rangle, \langle\psi|\rho|\psi\rangle > 0$
- Pure state  $\rho = |\psi\rangle\langle\psi|$  or  $\rho^2 = \rho$
- Mixed state  $\rho \neq |\psi\rangle\langle\psi|$
- Measured value of an observable, i.e. a Hermitian operator  $A$ , is  $\text{Tr}(\rho A)$
- Eigenvalues  $0 \leq \rho_j \leq 1, \sum_{j=1}^n \rho_j = 1$
- **Ex:** Check that the Bloch Sphere for a single qubit  $n = 2$

$$\rho(x, y, z) = \frac{1}{2} \begin{pmatrix} 1+z & x-iy \\ x+iy & 1-z \end{pmatrix}$$

satisfies all the requirements.

- Hilbert-Schmidt distance:

$$d_{HS}^2(\rho_A, \rho_B) = \text{Tr}((\rho_A - \rho_B)^2)$$

Joint probability density function for the eigenvalues

$$P(\rho_1, \dots, \rho_n) = \frac{1}{C_n} \mathbb{1}_{\sum_{j=1}^n \rho_j = 1} \prod_{1 \leq j < k \leq n} (\rho_k - \rho_j)^2$$

- Bures-distance [Bures 1969]

$$d_B^2(\rho_A, \rho_B) = 2 - 2\text{Tr} \sqrt{(\sqrt{\rho_A} \rho_B \sqrt{\rho_A})}$$

**Ex:** Joint probability density function for the eigenvalues [Hall 1998]

$$P(\rho_1, \dots, \rho_n) = \frac{1}{C_n} \mathbb{1}_{\sum_{j=1}^n \rho_j = 1} \prod_{j=1}^n \rho_j^{-1/2} \times \prod_{1 \leq j < k \leq n} \frac{(\rho_k - \rho_j)^2}{(\rho_k + \rho_j)}$$

- **Challenge Ex:** Show

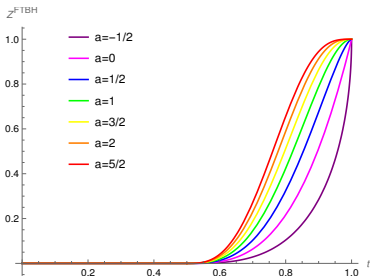
$$C_n = \frac{2^{-n(n-1)} \pi^{n/2}}{\Gamma(n^2/2)} \prod_{i=1}^n \Gamma(i+1)$$

## An example

Generalised gap probability for the Bures-Hall fixed-trace ensemble

$$Z^{\text{FT}}(t; n, a, \xi) := \frac{1}{C^{\text{FT}}(n, a, 1)} \frac{1}{n!} \left( \int_0^\infty -\xi \int_t^\infty \right) d\rho_1 \dots \left( \int_0^\infty -\xi \int_t^\infty \right) d\rho_n \\ \times \delta\left(\sum_{j=1}^n \rho_j - 1\right) \prod_{j=1}^n \rho_j^a \prod_{1 \leq j < k \leq n} \frac{(\rho_k - \rho_j)^2}{(\rho_k + \rho_j)}$$

Here  $t > 0$ ,  $\xi \in \mathbb{C}$ ,  $\text{Re}(a) > -1$ .



Fixed-trace Bures-Hall gap probability  $n = 2$ ,  $\xi = 1$  case of  $Z^{\text{FT}}(t; 2, a, 1)$

Step 1 :

Generalised gap probability for the Bures-Hall fixed-trace ensemble

$$\begin{aligned}
 Z^{\text{FT}}(t; n, a, r, \xi) &:= \frac{1}{C^{\text{FT}}(n, a, 1)} \frac{1}{n!} \left( \int_0^\infty -\xi \int_t^\infty \right) d\rho_1 \dots \left( \int_0^\infty -\xi \int_t^\infty \right) d\rho_n \\
 &\quad \times \delta\left(\sum_{j=1}^n \rho_j - r\right) \prod_{j=1}^n \rho_j^a \prod_{1 \leq j < k \leq n} \frac{(\rho_k - \rho_j)^2}{(\rho_k + \rho_j)}
 \end{aligned}$$

Here  $r > 0, t > 0, \xi \in \mathbb{C}, \text{Re}(a) > -1$ .

Step 2 :

Generalised gap probability for the unconstrained Bures-Hall ensemble

$$\begin{aligned}
 Z^{\text{U}}(s; n, a, \xi) &:= \frac{1}{C^{\text{U}}(n, a)} \frac{1}{n!} \left( \int_0^\infty -\xi \int_s^\infty \right) dx_1 \dots \left( \int_0^\infty -\xi \int_s^\infty \right) dx_n \\
 &\quad \prod_{1 \leq j < k \leq n} \frac{(x_k - x_j)^2}{(x_k + x_j)} \prod_{j=1}^n x_j^a e^{-x_j}
 \end{aligned}$$

Step 3 :

Generalised gap probability for the Cauchy-Laguerre Two Matrix Model

$$\begin{aligned}
 Z_n^{\text{CL}}(s, t; a, b; \xi, \psi) &:= \frac{1}{C^{\text{CL}}(n, a, b)} \frac{1}{(n!)^2} \left( \int_0^\infty -\xi \int_s^\infty \right) dx_1 \dots \left( \int_0^\infty -\xi \int_s^\infty \right) dx_n \\
 &\quad \left( \int_0^\infty -\psi \int_t^\infty \right) dy_1 \dots \left( \int_0^\infty -\psi \int_t^\infty \right) dy_n \\
 &\quad \times \prod_{j=1}^n x_j^a e^{-x_j} y_j^b e^{-y_j} \frac{\prod_{1 \leq j < k \leq n} (x_k - x_j)(y_k - y_j)}{\prod_{1 \leq j, k \leq n} (x_j + y_k)} \prod_{1 \leq j < k \leq n} (x_k - x_j)(y_k - y_j)
 \end{aligned}$$

Step 4 :

The Unconstrained Bures-Hall generating function is related to the Cauchy-Laguerre Two Matrix Model by

$$\left( Z^{\text{U}}(s; n, a, \xi) \right)^2 = 2^n Z_n^{\text{CL}}(s, s; a, a + 1, \xi, \xi)$$

extending [Forrester+Kieburg 2016]

Bi-moments  $M_{j,k}$  are defined by

$$M_{j,k} = \int_{S_1 \times S_2} dx dy \frac{w(x,y)}{x+y} x^j y^k < \infty, \quad \forall j, k \geq 0.$$

The Cauchy-Laguerre gap probability is a determinant of the truncated Gram matrix

$$Z_n^{\text{CL}} = \det(M_{j,k})_{j,k=0}^{n-1} \neq 0, \quad n \geq 1$$

with

$$w(x,y) = w_1(x)w_2(y) = x^a e^{-x} y^b e^{-y}, \quad \text{Re}(a,b) > -1$$

## Proposition (Bertola+Gekhtman+Szmigielski 2010)

Let two sequences of univariate moments be defined as

$$\alpha_j := \int_{S_1} dx w_1(x) x^j, \quad \beta_j := \int_{S_2} dy w_2(y) y^j.$$

A key identity, which we call the Cauchy relation because it follows directly from the Cauchy kernel, for the bi-moments is the rank-1 condition

$$M_{j+1,k} + M_{j,k+1} = \alpha_j \beta_k.$$



Define an *inner product* over polynomial spaces  $\cup_{n \geq 0} \Pi_n[x]$  using

$$\langle f, g \rangle := \int_{S_1 \times S_2} dx dy \frac{w(x, y)}{x + y} f(x)g(y)$$

with  $f, g \in \cup_{n \geq 0} \Pi_n[x]$

Two sequences of *bi-orthogonal polynomials*  $\{P_n(x), Q_n(y)\}_{n=0}^{\infty}$  satisfying the orthogonality relation

$$\langle P_m, Q_n \rangle = \delta_{m,n}$$

Let two sequences of univariate moments be defined as

$$\pi_j := \int_{S_1} dx w_1(x) P_j, \quad \eta_j := \int_{S_2} dy w_2(y) Q_j.$$

## Proposition (BGS 2010)

Let the coefficients  $\pi_n, \eta_n$  be non-vanishing for all  $n \geq 0$ , and the norms  $h_n$  or  $Z_n^C$  similarly be non-vanishing. The bi-orthogonal polynomials  $P_n(x), Q_n(y)$  defined by the general orthogonality condition satisfy uncoupled, third-order scalar recurrence relations of the form

$$x \left( \frac{1}{\pi_{n+1}} P_{n+1} - \frac{1}{\pi_n} P_n \right) = r_{n,2} P_{n+2} + r_{n,1} P_{n+1} + r_{n,0} P_n + r_{n,-1} P_{n-1},$$

$$y \left( \frac{1}{\eta_{n+1}} Q_{n+1} - \frac{1}{\eta_n} Q_n \right) = s_{n,2} Q_{n+2} + s_{n,1} Q_{n+1} + s_{n,0} Q_n + s_{n,-1} Q_{n-1}.$$

Define the reproducing kernel

$$K_n^{0,0}(x, y) := \sum_{l=0}^n P_l(x)Q_l(y).$$

In [BGS 2010] the notion of a pair of *intertwined* polynomials was introduced

$$\hat{P}_n(x) := - \sum_{k=0}^n \eta_k P_k(x),$$

$$\check{Q}_n(y) := - \sum_{j=0}^n \pi_j Q_j(y),$$

### Proposition (WW 2022)

Let all the standard conditions apply. The Christoffel-Darboux sum for the reproducing kernel has the evaluation

$$(x + y)K_n^{0,0}(x, y) = \hat{P}_n(x)\check{Q}_n(y) + \frac{S_n}{S_{n+1}} [P_n(x)Q_{n+1}(y) + P_{n+1}(x)Q_n(y)].$$

*Associated Functions of the first type of  $n$ -th order*

$$P_n^{(1)}(z) := \int_{S_1} dx \frac{w_1(x)}{z-x} P_n(x), \quad Q_n^{(1)}(z) := \int_{S_2} dy \frac{w_2(y)}{z-y} Q_n(y)$$

*Associated Functions of the second type*

$$P_n^{(2)}(z) := \int_{S_1 \times S_2} dx dy \frac{w(x,y)}{x+y} \frac{P_n(x)}{z-x}, \quad Q_n^{(2)}(z) := \int_{S_1 \times S_2} dx dy \frac{w(x,y)}{x+y} \frac{Q_n(y)}{z-y}$$

$3 \times 3$  matrix variables for  $n \geq 0$

$$\mathcal{P}_n := \begin{pmatrix} P_{n+1}(x) & P_{n+1}^{(1)}(x) & P_{n+1}^{(2)}(x) \\ P_n(x) & P_n^{(1)}(x) & P_n^{(2)}(x) \\ P_{n-1}(x) & P_{n-1}^{(1)}(x) & P_{n-1}^{(2)}(x) \end{pmatrix}$$

Matrix system of first order recurrences with transfer matrices  $K_n$  and  $L_n$

$$\mathcal{P}_{n+1} = K_n \mathcal{P}_n, \quad \mathcal{Q}_{n+1} = L_n \mathcal{Q}_n$$

## Proposition (WW 2022)

The kernels  $K_n^{\mu,\nu}(x, y)$  for  $\mu, \nu = 0, 1, 2$  are the bilinear forms

$$\pi_n \eta_n(x + y) K_n^{\mu,\nu}(x, y) = \mathcal{P}_n^{(\mu)}(x)^T G_n(x, y) \mathcal{Q}_n^{(\nu)}(y)$$

where  $G(x, y)$  is the  $3 \times 3$  matrix

$$G_n(x, y) = \begin{pmatrix} \frac{S_n^2}{S_{n+1}^2} & \frac{S_n}{S_{n+1}} [Y_{n,n} - y] & -\frac{S_{n-1}}{S_{n+1}} \\ \frac{S_n}{S_{n+1}} [X_{n,n} - x] & [X_{n,n} + y] [Y_{n,n} + x] & \frac{S_{n-1}}{S_n} [Y_{n,n} + x] \\ -\frac{S_{n-1}}{S_{n+1}} & \frac{S_{n-1}}{S_n} [X_{n,n} + y] & \frac{S_{n-1}^2}{S_n^2} \end{pmatrix}$$



Deduce three *Schlesinger equations*

$$\partial_s \mathcal{A}_n - \partial_x \mathcal{B}_n = [\mathcal{B}_n, \mathcal{A}_n], \quad \partial_t \mathcal{A}_n - \partial_x \mathcal{C}_n = [\mathcal{C}_n, \mathcal{A}_n], \quad \partial_t \mathcal{B}_n - \partial_s \mathcal{C}_n = [\mathcal{C}_n, \mathcal{B}_n]$$

$\sigma$ -functions

$$\sigma_n(s, t) := s \partial_s \log Z_n^{\text{CL}}, \quad \tau_n(s, t) := t \partial_t \log Z_n^{\text{CL}}$$

The  $\sigma$ -functions have evaluations in terms of *reproducing kernels*

$$\sigma_n(s, t) = -\xi s^{a+1} e^{-s} K_{n-1}^{0,1}(s, -s), \quad \tau_n(s, t) = -\psi t^{b+1} e^{-t} K_{n-1}^{1,0}(-t, t)$$

*Kernels* defined by

$$K_n^{0,1}(x, y) := \sum_{l=0}^n P_l(x) Q_l^{(1)}(y), \quad K_n^{1,0}(x, y) := \sum_{l=0}^n P_l^{(1)}(x) Q_l(y)$$

The  $\sigma$ -forms for deformed Cauchy-Laguerre bi-orthogonal ensemble have the evaluations

$$\sigma_{n+1}(s, t) = s \text{Tr} \begin{pmatrix} 0 & \frac{\pi_{n+1}}{\pi_n} & 0 \\ 0 & 1 & 0 \\ 0 & \frac{\pi_{n-1}}{\pi_n} + \frac{S_n}{S_{n-1}} & 0 \end{pmatrix} \mathcal{A}_n^{(s)} + \text{Tr} (\mathcal{A}_n^{(\Sigma)} \mathcal{A}_n^{(s)}) - \frac{t}{(s+t)} \text{Tr} (\mathcal{A}_n^{(-t)} \mathcal{A}_n^{(s)}),$$

A number of statistics have been calculated for the Bures-Hall ensemble which are global in nature such as:

- Quantum purity  $\mathbb{E}[\text{Tr}(\rho^2)]$ .
- von Neumann entropy  $\mathbb{E}[-\text{Tr}(\rho \log \rho)]$  and its higher moments.

However

- The gap probability is a refined local statistic of the density matrix spectrum, especially of the extreme eigenvalues.
- It is now accessible with the tools of "integrable probability" and Riemann-Hilbert methods.
- It is a rare example of an integrable formulation of a *Pfaffian point processes*, whereas virtually all the other known cases are of *determinantal point processes*.
- It is a rank 3 isomonodromic system and is beyond the Painlevé class which has rank 2.

# The Cast



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